

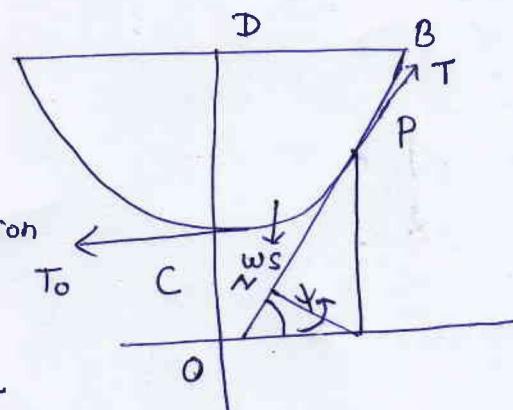
**Model Answer**  
**B. Sc. 5<sup>th</sup> Semester Mathematics**  
**Department of Pure & Applied Mathematics**  
**GGV, Bilaspur**

**Subject: Mechanics-I**

**Paper Code: AS-2829**

1.(a). Let  $ACB$  be the string of which  $C$  is the lowest point.

$P$  is any point of the string such that the length of the arc  $CP = s$ . Let  $T$  be the tension at  $P$  and  $\gamma$  its inclination with horizontal.



At the lowest point  $C$  tension is horizontal and let it be  $T_0$ .

Let  $w$  be weight per unit length. Then the portion  $CP$  of the string is in equilibrium under the action of the forces

(i) The tension  $T$  at  $P$  and  $T_0$  at  $C$ .

(ii) The weight  $ws$  vertically downward

Resolving these forces horizontally and vertically, we have

$$T \cos \gamma = T_0 \quad (1)$$

$$T \sin \gamma = ws \quad (2)$$

Let the tension at lowest point  $T_0$  to be equal to the weight of the string of length  $c$  then

$$T_0 = wc \quad (3)$$

dividing (2) by (1) and putting value of  $T_0$ , we get

$$\tan \gamma = s/c$$

$$\Rightarrow \boxed{s = c \tan \gamma}$$

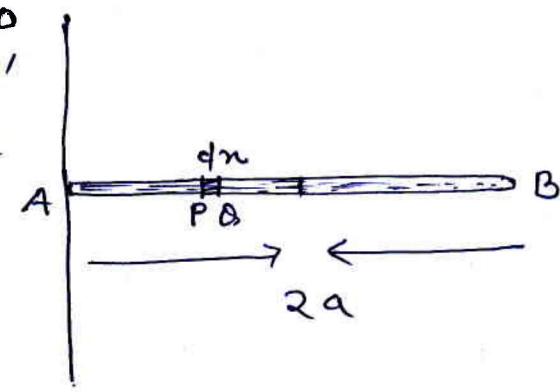
(b). Kinetically Equivalent system: Two mechanical systems that are such that their moments of inertia about all lines are the same, are called equimomental or kinetically equivalent systems.

Ans.

- (c). Let AB be the rod and  $OAO'$  is the axis passing through one vertex A and perpendicular to the rod.

$$\text{Let } AP = n \text{ and } PQ = dn,$$

$$\text{the mass of element } PQ = \frac{M}{2a} dn.$$



Now moment of inertia of  $PQ$  about  $OAO'$

$$= \left( \frac{M}{2a} dn \right) n^2$$

Therefore moment of inertia of whole rod about  $OAO'$

$$= \int_0^{2a} \frac{M}{2a} n^2 dn$$

$$= \frac{M}{2a} \left[ \frac{n^3}{3} \right]_0^{2a}$$

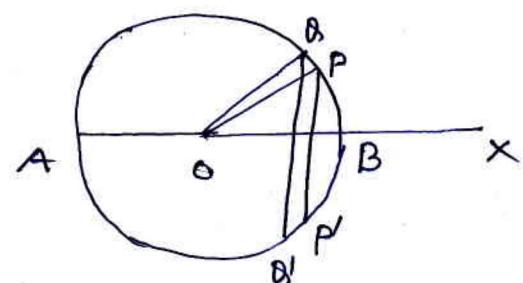
$$= \frac{M}{2a} \times \frac{8a^3}{3} = \frac{4Ma^2}{3}$$

Ans.

- (d). Let the surface of the sphere be divided into circular bands ~~like~~ like  $PQ'P'Q'$  perpendicular to  $OA$ .

$$\text{if } \angle POn = \theta$$

$$\text{and } \angle QOP = d\theta$$



(3)

$$\text{The surface of the band} = 2\pi a \sin \theta \cdot a d\theta$$

Now moment of inertia about  $o$  of the band

$$= \frac{M}{4\pi a^2} (2\pi a \sin \theta \cdot a d\theta) a^2 \sin^2 \theta$$

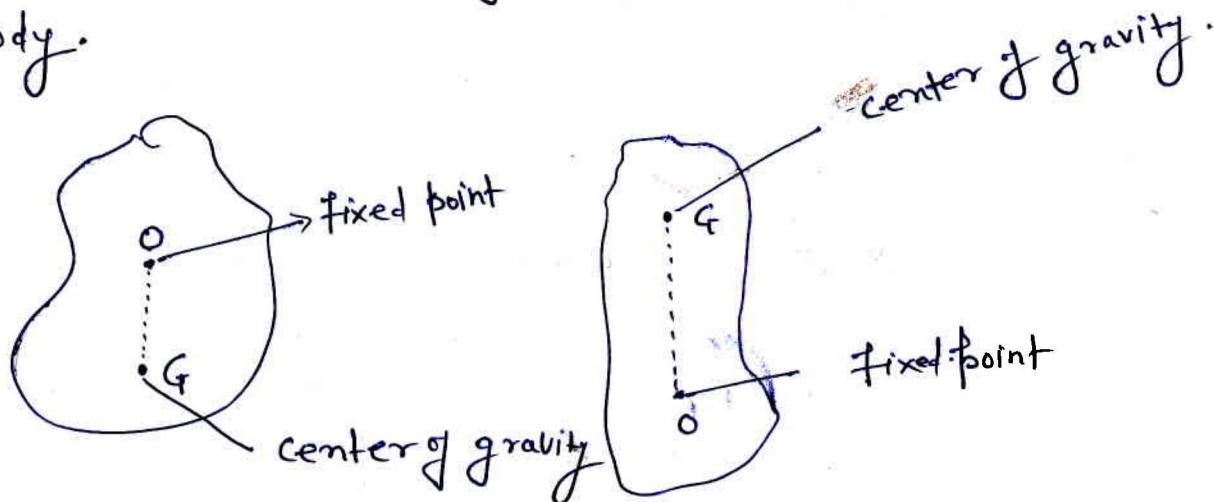
Hence, moment of inertia of the follow sphere about  $o$

$$= \frac{M \cdot a^2}{2} \int_0^{\pi} \sin^3 \theta d\theta$$

$$= M \frac{2a^2}{3}$$

Ans.

(e). Equilibrium of a rigid body: If a rigid body be in equilibrium, one point only of the body being fixed and the center of gravity of the body will be on the vertical line passing through the fixed point of the body.



These two bodies are said to be in state of stable and unstable equilibrium.

(4)

Energy test for stability: We know that the sum of kinetic and potential energy is constant. In the position of equilibrium kinetic energy is zero, therefore potential energy is constant i.e. potential energy is either maximum or minimum. Now if the system be slightly displaced from a position of maximum potential energy and if the potential energy decreases i.e. K.E. increases, or in other words kinetic energy will be positive as in the equilibrium position it was zero. Hence the system moves further away from the position of maximum potential energy showing that the equilibrium in the position of maximum potential energy is an unstable one. Further in the case of minimum potential energy we have stable equilibrium.

Maximum potential energy  $\Rightarrow$  Unstable equilibrium

Minimum potential energy  $\Rightarrow$  stable equilibrium.

(5)

For Common Catenary, We know that

$$y = c \sec \gamma \quad \text{--- (I)}$$

and

$$s = c \tan \gamma \quad \text{--- (II)}$$

From (I) and (II), We can write

$$\begin{aligned} y^2 - s^2 &= c^2 (\sec^2 \gamma - \tan^2 \gamma) \\ &= c^2 \end{aligned}$$

$$\Rightarrow y^2 = c^2 + s^2$$

Ans.

Product of inertia: The product of inertia of a lamina with respect to the axes  $ox$  and  $oy$  is defined by  $\int xy \, dm$ ,  $dm$  is elementary mass.

Since for any curve, we know that

$$\frac{dn}{ds} = \cos \gamma \quad \text{--- (I)}$$

$$\frac{dy}{ds} = \sin \gamma \quad \text{--- (II)}$$

Also we can write

$$\frac{dn}{d\gamma} = \frac{dn}{ds} \cdot \frac{ds}{d\gamma}$$

$$\begin{aligned} \Rightarrow \frac{dn}{d\gamma} &= \cos \gamma \frac{ds}{d\gamma} (c \tan \gamma), \quad \text{since for catenary } \\ &\qquad\qquad\qquad s = c \tan \gamma \\ &= c \cos \gamma \cdot \sec^2 \gamma \\ \Rightarrow dn &= c \sec \gamma \, d\gamma \end{aligned}$$

Integrating both sides, we have

$$n = c \log(\sec y + \tan y) + A \quad \text{--- (III)}$$

Since at  $n=0, y=0$ , From (III), we have

$A=0$ , putting in (III), we get

$$\boxed{n = c \log(\sec y + \tan y)}$$

Ans

2(b). We know that the equation of common catenary is

$$\begin{aligned} y &= c \cosh \frac{x}{c} \\ &= \frac{c}{2} \left( e^{x/c} + e^{-x/c} \right) \\ &= c \left[ 1 + \frac{1}{2!} \left( \frac{x}{c} \right)^2 + \frac{1}{4!} \left( \frac{x}{c} \right)^4 + \dots \right]. \end{aligned} \quad \text{--- (I)}$$

Since span ( $x$ ) is very small and  $c$  is very large, therefore,  $\frac{x}{c}$  will be very small and the powers of  $\frac{x}{c}$  beyond two can be neglected, so that equation

(I) can be written as

$$y = c \left( 1 + \frac{x^2}{2c^2} \right)$$

$$\Rightarrow x^2 = 2c(y - c) \quad \text{--- (II)}$$

Equation (II) represents a parabola with latus rectum  $2c$  or  $\frac{2T_0}{\omega}$ , where  $\omega$  is weight per unit length and  $T_0 (= \omega c)$  is tension at the lowest point.

Ans.

(3.).

7 8

ACB is the chain of length  $l$ .

AB is breadth of the river,

and  $CD = R$  (given)

We know that

$$y = c \cosh \frac{x}{c} \quad (i)$$

At the point B,  $y = c+R$ , therefore we have

$$c+R = c \left[ 1 + \frac{n^2}{2c^2} + \frac{1}{3!} \frac{n^3}{c^3} + \dots \right]$$

$$c+R = c + \frac{n^2}{2c} \quad \left( \frac{n}{c} \text{ is very small neglecting higher terms} \right)$$

$$\Rightarrow R = \frac{n^2}{2c}$$

$$\Rightarrow c = \frac{n^2}{2R} \quad (ii)$$

Now we know that

$$x = c \sinh \frac{x}{c}$$

$$\frac{l}{2} = c \left[ \frac{n}{c} + \frac{n^3}{3! c^3} + \dots \right]$$

$$\Rightarrow x \frac{l}{2} = 2n + \frac{n^3}{8c^2} \times x$$

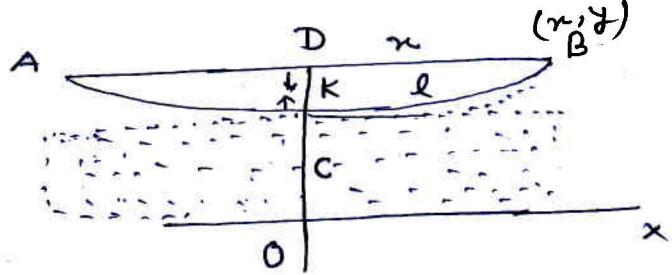
$$\Rightarrow l - 2n = \frac{n^3}{3c^2} = \frac{n^3}{3} \frac{8K^2}{n^4} \quad \left[ \text{from (ii)} \right]$$

$$\Rightarrow l - 2n = \frac{8K^2}{3n} = \frac{8 \cdot K^2}{3l} \quad \text{since } n \approx l$$

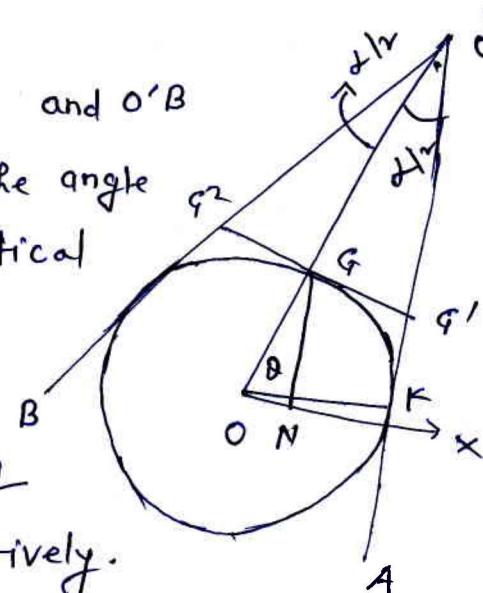
$\Rightarrow$  Length of the Chain - Breadth of the river

$$= \frac{8}{3} \frac{K^2}{l}$$

Ans.



(4). Two equal uniform rods  $O'A$  and  $O'B$  are joined at  $O'$  so that the angle between them is  $\alpha$  in vertical plane on a smooth sphere of radius  $r$ .  $G'$  and  $G^2$  are center of gravity of the rods  $O'A$  and  $O'B$  respectively.



Let  $G$  be their combined C.G. then  $OG$  will bisect the angle  $BO'A$  and also the base  $G'G^2$ .

$$\text{Let } O'A = O'B = 2a$$

Now the height of C.G. of the system above the horizontal

$$\begin{aligned} z &= GN = (OG \sin \theta) - \\ &= (OO' - O'G) \sin \theta \quad \text{--- (i)} \end{aligned}$$

Now from right angle triangle  $O'OK$  and  $O'GG_1$  we have

$$\sin \frac{\alpha}{2} = \frac{r}{OO'}, \quad \cos \frac{\alpha}{2} = \frac{O'G}{O'G_1}$$

$$\Rightarrow OO' = r \operatorname{cosec} \frac{\alpha}{2}, \quad O'G = a \cos \frac{\alpha}{2}$$

Putting these values in (i), we get

$$z = \left( r \operatorname{cosec} \frac{\alpha}{2} - a \cos \frac{\alpha}{2} \right) \sin \theta$$

$$\Rightarrow \frac{dz}{d\theta} = \left( r \operatorname{cosec} \frac{\alpha}{2} - a \cos \frac{\alpha}{2} \right) \cos \theta$$

$$\text{and } \frac{d^2 z}{d\theta^2} = \left( a \cos \frac{\alpha}{2} - 2 \cosec \frac{\alpha}{2} \right) \sin \theta \quad (9)$$

For position of equilibrium,

$$\begin{aligned} \frac{dz}{d\theta} = 0 &\Rightarrow \left( a \cos \frac{\alpha}{2} - 2 \cosec \frac{\alpha}{2} \right) \cos \theta = 0 \\ &\Rightarrow \cos \theta = 0 \\ &\Rightarrow \theta = \pi/2 \end{aligned}$$

$$\text{Now } \left( \frac{d^2 z}{d\theta^2} \right)_{\theta=\frac{\pi}{2}} = \left( a \cos \frac{\alpha}{2} - 2 \cosec \frac{\alpha}{2} \right)$$

Hence the equilibrium will be stable or unstable according as  $z$  is Min. or Max. i.e.

$$a \cos \frac{\alpha}{2} > \text{ or } < 2 \cosec \frac{\alpha}{2}$$

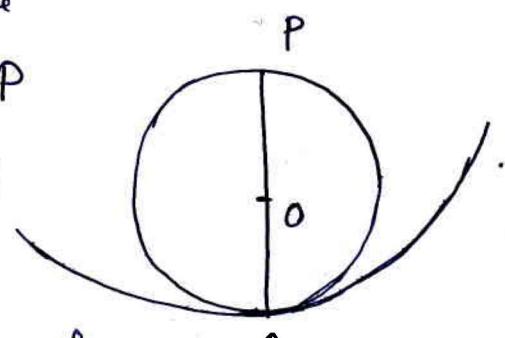
$$\text{or } 2a \cos \frac{\alpha}{2} < \text{ or } < 2 \cosec \frac{\alpha}{2}$$

$$\text{or } 2a <$$

$$\Rightarrow 2a > \text{ or } < 4 \cosec \alpha .$$

Ans

- (5). Let  $W$  be the weight of the sphere of radius  $r$  and let  $P$  be the weight attached to its lowest highest point. If  $h$  be the height of the C.G. of



the system above  $A$ , the point of contact of the sphere and the bowl,

$$f = \frac{W \cdot \gamma + P \cdot 2\gamma}{W+P}$$

(10)

For stable equilibrium

$$\frac{1}{f} > \frac{1}{\gamma} - \frac{1}{R}$$

where  $R$  is the radius of the bowl (here  $\gamma$  and  $R$  are measured in opposite directions).

$$\therefore \frac{W+P}{\gamma(W+P)} > \frac{1}{\gamma} - \frac{1}{2\gamma}, \text{ since } R = \gamma$$

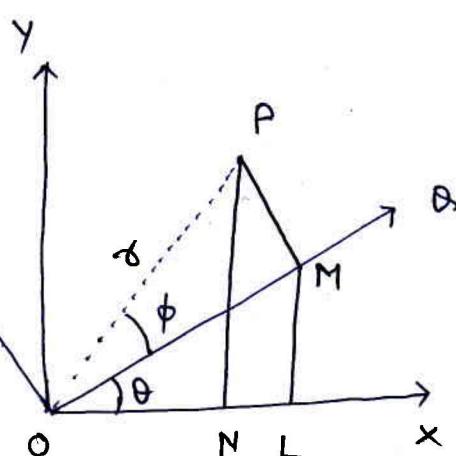
$$\text{or } W+P > \frac{1}{2}(W+2P)$$

$$\text{or } W > \frac{W}{2}$$

Which is always true and independent of  $P$ . Hence the equilibrium is stable for all values of  $P$ .

(6).

Let  $m$  be the element  $dm$  of mass of the lamina enclosing a point  $P(x, y)$  referred to  $ox, oy$  as axes.



Let  $OQ$  is any line making an angle  $\theta$  with  $ox$ .

Let  $A, B$  be the moments of inertia of the lamina about  $ox$  and  $oy$  and  $F$ , its product of

(1)

inertia about  $ox, oy$ . Then, if  $PM$  is the perpendicular from  $P$  on  $OQ$ , moment of inertia of the lamina about  $OQ$

$$\begin{aligned}
 &= \sum m (PM)^2 \\
 &= \sum m (y \cos \theta - x \sin \theta)^2 \\
 &= \sum \cos^2 \theta (my^2) + \sum \sin^2 \theta (mx^2) \\
 &\quad - 2 \sin \theta \cos \theta \sum mxny \\
 &= \cos^2 \theta \sum my^2 + \sin^2 \theta \sum mx^2 \\
 &\quad - 2 \sin \theta \cos \theta \sum mxny \\
 &= A \cos^2 \theta + B \sin^2 \theta - 2F \sin \theta \cos \theta
 \end{aligned}$$

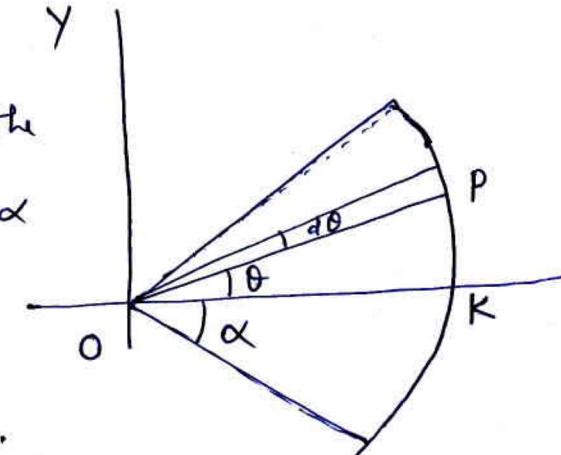
Now the product of inertia about  $OQ$  and  $OR$

$$\begin{aligned}
 &= \sum m OM \cdot PM \\
 &= \sum m (x \cos \theta + y \sin \theta) (y \cos \theta - x \sin \theta) \\
 &= \sin \theta \cos \theta \sum my^2 - \sin \theta \cos \theta \sum mx^2 \\
 &\quad + (\cos^2 \theta - \sin^2 \theta) \sum mxny \\
 &= (A - B) \sin \theta \cos \theta + F \cos 2\theta
 \end{aligned}$$

Ans

7. Let  $O$  be the centre  $y$

and  $K$  the middle point of the arc subtending an angle  $2\alpha$  at  $O$ . If  $M$  be the mass of an element the arc of the circle of radius  $a$ .



$$\text{Mass per unit length} = \frac{M}{2a\alpha}$$

$$\text{Mass of an element } d\theta \text{ at } P = \frac{M}{2a\alpha} d\theta$$

$$\text{Hence moment of inertia about } OK = \int_{-a}^a \frac{M}{2a\alpha} d\theta a^2 \sin^2\theta$$

$$\text{and moment of inertia about } Oy = \int_{-a}^a \frac{M}{2a\alpha} d\theta a^2 \cos^2\theta$$

therefore moment of inertia about an axis through O perpendicular to the plane = sum of moments of inertia about OK and about Oy

$$= \int_{-a}^a \frac{M}{2a\alpha} (a^2 \sin^2\theta + a^2 \cos^2\theta) d\theta$$

$$= M a^2$$

Ans

Q- (a). Compound Pendulum: A compound pendulum is one where the rod is not massless, and may have extended size, that is, an arbitrarily shaped rigid body swinging by a pivot. In this case pendulum periods depends on its moment of inertia  $I$  around the pivot point.

The equation of torque is given by

$$\tau = I\alpha,$$

where  $\alpha$  is the angular acceleration,  $\tau$  is the torque. ~~.....~~

(13)

Since torque is generated by gravity so,

$$\tau = -mgL \sin\theta$$

Where  $L$  is the distance from the pivot to the centre of mass of pendulum.  $\theta$  is the angle from the vertical.

For small angle approximation  $\sin\theta \approx \theta$

$$\alpha \approx \frac{mgL\theta}{I}$$

This gives a period of

$$T = 2\pi \sqrt{\frac{I}{mgL}}$$

Ans

(b). D'Alembert Principle: The principle states that the sum of the differences between the forces acting on a system of mass particles and time derivatives of the momenta of the system itself along any virtual displacement consistent with the constraints of the system, is zero. Symbolically we can write

as

$$\sum_i (F_i - m_i \cdot a_i) \cdot \delta r_i = 0,$$

Where

$i$  is an integer

$F_i$  is the total applied force on  $i$ th particle

$m_i$  is mass of  $i$ th particle

$a_i$  is acceleration of  $i$ th particle

$\delta r_i$  is virtual displacement of the  $i$ th particle

This principle is also known as Lagrange's D'Alembert principle.  
It is named after its discoverer, the French physicist  
and Mathematician Jean le Rond d'Alembert.



(Dr. B. B. Chaturvedi)  
Department of Pure & Applied Mathematics  
GGV, Bilaspur